

Thermal expansion, elastic stress and finite deformation kinematics

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ABSTRACT

Kinematics + thermal expansion \implies

\implies temperature gradient always causes elastic stress

Standard kinematics:

reference frame, reference t_0 , reference state/configuration:

material point $P \longleftrightarrow$ its position \mathbf{X} at t_0

at t , its position is $\mathbf{x} = \mathbf{X} + \mathbf{u}_{t_0 \rightarrow t}(\mathbf{X})$

For small strains:

Cauchy strain: $\mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} = (\mathbf{u}_{t_0 \rightarrow t} \otimes \nabla)^S \quad [\nabla \equiv \overleftarrow{\nabla} \equiv \overrightarrow{\nabla}]$

with the property

$$\nabla \times \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} \times \nabla = \mathbf{0}$$

(Saint-Venant compatibility condition)

For nonsmall strains:

Deformation gradient: $\mathbf{F}_{t_0 \rightarrow t} = \mathbf{1} + \mathbf{u}_{t_0 \rightarrow t} \otimes \nabla$

Left Cauchy-Green = Finger deformation tensor:

$$\mathbf{B}_{t_0 \rightarrow t} = \mathbf{F}_{t_0 \rightarrow t} \mathbf{F}_{t_0 \rightarrow t}^{\top}$$

Strain tensors: for any number n ,

$$\mathbf{E}_{t_0 \rightarrow t}^{(n)} = \frac{1}{n} \left(\mathbf{B}_{t_0 \rightarrow t}^{n/2} - \mathbf{1} \right), \quad \mathbf{E}_{t_0 \rightarrow t}^{(0)} = \ln \mathbf{B}_{t_0 \rightarrow t}^{1/2}$$

Elasticity:

$$\boldsymbol{\sigma} = 2G \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} + \left(K - \frac{2}{3}G \right) \left(\text{tr} \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} \right) \mathbf{1},$$

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The recent approach to kinematics: (Fülöp–Ván, 2010)

Material point: \mathbf{X} at $t_0 \longrightarrow P \in \mathcal{C}$ (material manifold)

Motion: $\mathbf{u}_{t_0 \rightarrow t}(\mathbf{X}) \longrightarrow$ at t , P is at spacetime point $r_t(P)$

Local geometry of the motion:

$$\mathbf{F}_{t_0 \rightarrow t} = \mathbf{1} + \mathbf{u}_{t_0 \rightarrow t} \otimes \nabla \longrightarrow \mathbf{J}_t(P) = (r_t \otimes \tilde{\nabla})(P)$$

Current metric (flat): $?$ $\longrightarrow d_t(P, Q) = \|r_t(Q) - r_t(P)\|_{\mathbf{h}},$

$$\tilde{\mathbf{h}}_t = \mathbf{J}_t^\top \mathbf{h} \mathbf{J}_t.$$

Solids: relaxed/natural metric: $?$ $\longrightarrow \tilde{\mathbf{g}}$

Kinematic quantity for elasticity:

How far is the current metric from the natural metric:

Elastic shape: $\tilde{\mathbf{A}}_t = \tilde{\mathbf{g}}^{-1} \tilde{\mathbf{h}}_t$, $\mathbf{A}_t = \mathbf{J}_t \tilde{\mathbf{A}}_t \mathbf{J}_t^{-1}$

Deformedness: $\mathbf{D}_t = \ln \mathbf{A}_t^{1/2}$

In the natural/relaxed state: $\tilde{\mathbf{h}}_t = \tilde{\mathbf{g}}$, $\mathbf{A}_t = \mathbf{1}$, $\mathbf{D}_t = \mathbf{0}$

Elastic constitutive relation:

$$\boldsymbol{\sigma}_t = 2G\mathbf{D}_t + \left(K - \frac{2}{3}G\right) (\text{tr}\mathbf{D}_t) \mathbf{1}$$

The possibility for $\tilde{\mathbf{h}}_t = \tilde{\mathbf{g}} \implies \tilde{\mathbf{g}}$ also flat: its $\mathbf{R}^{\text{Ricci}} = \mathbf{0}$:

$$\begin{aligned}
& \text{tr}_{1,5;3,4} [\mathbf{A}^{-1} \otimes \mathbf{A}^{-1} (\mathbf{A} \otimes \nabla) (\mathbf{A} \otimes \nabla)] + \\
& \quad + 2\mathbf{A}^{-1} \text{tr}_{1,2} [\mathbf{A} (\nabla \otimes \nabla \otimes \mathbf{A})] \mathbf{A}^{-1} + \\
& \quad + 2\text{tr}_{2,3} [\mathbf{A}^{-1} \otimes (\nabla \cdot \mathbf{A}) (\nabla \otimes \mathbf{A})] \mathbf{A}^{-1} + \\
& \quad + 2\text{tr}_{1,2} [\mathbf{A}^{-1} (\mathbf{A} \otimes \nabla \otimes \nabla)] - 2\mathbf{A}^{-1} \text{tr}_{2,4} [(\mathbf{A} \otimes \nabla \otimes \nabla)] + \\
& \quad + 2\text{tr}_{1,4;3,5} [\mathbf{A}^{-1} \otimes \mathbf{A}^{-1} (\mathbf{A} \otimes \nabla) (\mathbf{A} \otimes \nabla)] + \\
& \quad + \text{tr}_{1,2;3,5} [\mathbf{A}^{-1} (\mathbf{A} \otimes \nabla) (\mathbf{A} \otimes \nabla) \otimes \mathbf{A}^{-1}] - \\
& \quad - 2\mathbf{A}^{-1} \text{tr}_{2,5;3,6} [(\mathbf{A} \otimes \nabla) \mathbf{A} (\nabla \otimes \mathbf{A}) \otimes \mathbf{A}^{-1}] \mathbf{A}^{-1} - \\
& \quad - 2 [\nabla \otimes (\nabla \cdot \mathbf{A})] \mathbf{A}^{-1} + 2\text{tr}_{2,4} [(\nabla \otimes \mathbf{A}) \mathbf{A}^{-1} (\mathbf{A} \otimes \nabla)] \mathbf{A}^{-1} - \\
& \quad - 3\text{tr}_{2,6;3,5} [(\nabla \otimes \mathbf{A}) \mathbf{A}^{-1} (\mathbf{A} \otimes \nabla) \otimes \mathbf{A}^{-1}] - \\
& \quad - \text{tr}_{3,4;2,5} [\mathbf{A}^{-1} \otimes \mathbf{A}^{-1} (\mathbf{A} \otimes \nabla) \mathbf{A} (\nabla \otimes \mathbf{A})] \mathbf{A}^{-1} - \\
& \quad - 2\mathbf{A}^{-1} \text{tr}_{2,4;3,5} [(\mathbf{A} \otimes \nabla) \otimes (\nabla \otimes \mathbf{A})] \mathbf{A}^{-1} = \mathbf{0}
\end{aligned}$$

Small deformedness: $\mathbf{D} \ll \mathbf{1}$:

$$\nabla \times \mathbf{D} \times \nabla + (\text{higher order terms}) = \mathbf{0}$$

$$(\text{recall } \nabla \times \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} \times \nabla = \mathbf{0})$$

Connection to the conventional kinematics:

$$\mathbf{F}_{t_0 \rightarrow t} = \mathbf{J}_t \mathbf{J}_{t_0}^{-1}$$

The initial state is not relaxed, in general;

$$\mathbf{A}_t = \mathbf{F}_{t_0 \rightarrow t} \mathbf{A}_{t_0} \mathbf{F}_{t_0 \rightarrow t}^\top \quad (\text{compare to } \mathbf{B}_{t_0 \rightarrow t} = \mathbf{F}_{t_0 \rightarrow t} \mathbf{F}_{t_0 \rightarrow t}^\top)$$

Thermal expansion:

$$\tilde{\mathbf{g}}^{(T)} = \Lambda^{(T_0, T)} \tilde{\mathbf{g}}^{(T_0)} \quad \Lambda^{(T_0, T)} : \text{thermal scaling factor}$$

If $T \approx T_0$:

$$\Lambda^{(T_0, T)} \approx [1 + \alpha^{(T_0)}(T - T_0)]^2 \approx 1 + 2\alpha^{(T_0)}(T - T_0)$$

$\alpha^{(T_0)}$: linear thermal expansion coefficient at T_0

If T_0 at t_0 , T at t :

$$\mathbf{A}_t = \frac{1}{\Lambda^{(T_0, T)}} \mathbf{F}_{t_0 \rightarrow t} \mathbf{A}_{t_0} \mathbf{F}_{t_0 \rightarrow t}^\top$$

If $T \approx T_0$, $\mathbf{D}_t \ll \mathbf{1}$, $\mathbf{F}_{t_0 \rightarrow t} \approx \mathbf{1}$:

$$\begin{aligned}\mathbf{D}_t &\approx \frac{1}{2} (\mathbf{A}_t - \mathbf{1}) \approx \frac{1}{2} (\mathbf{F}_{t_0 \rightarrow t} \mathbf{F}_{t_0 \rightarrow t}^\top - \mathbf{1}) - \frac{1}{2} (\Lambda^{(T_0, T)} - \mathbf{1}) \mathbf{1} \\ &\approx \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} - \alpha^{(T_0)} (T - T_0) \mathbf{1} ,\end{aligned}$$

and $\boldsymbol{\sigma}_t = 2G\mathbf{D}_t + \left(K - \frac{2}{3}G\right) (\text{tr}\mathbf{D}_t) \mathbf{1}$ expands to

$$\boldsymbol{\sigma}_t \approx 2G\mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} + \left(K - \frac{2}{3}G\right) \mathbf{E}_{t_0 \rightarrow t}^{\text{Cauchy}} - 3K\alpha^{(T_0)} (T - T_0) \mathbf{1} :$$

the Duhamel-Neumann expression

A general result: $\Lambda^{(T_0, T)}$ makes $\tilde{\mathbf{g}}$ space-and-time dependent
 $\implies \tilde{\mathbf{g}}$ is no longer flat
 $\implies \tilde{\mathbf{g}}$ cannot be equal to any $\tilde{\mathbf{h}}_t$

$$\begin{aligned} \mathbf{R}^{\text{Ricci}} &= \mathbf{R}^{\text{Ricci}} \Big|_{\Lambda=1} + \\ &+ \frac{1}{4\Lambda^2} \left\{ 3(\nabla\Lambda) \otimes (\nabla\Lambda) - 2\Lambda(\nabla \otimes \nabla\Lambda) - \Lambda \text{tr}_{2,3} [(\nabla \otimes \mathbf{A})(\nabla\Lambda)] \mathbf{A}^{-1} \right. \\ &- \Lambda \mathbf{A}^{-1} \text{tr}_{2,4} [(\mathbf{A} \otimes \nabla) \otimes (\nabla\Lambda)] - 2\Lambda [(\nabla\Lambda)(\mathbf{A} \cdot \nabla)] \mathbf{A}^{-1} \\ &- 2\Lambda \text{tr} [(\nabla \otimes \nabla\Lambda) \mathbf{A}] \mathbf{A}^{-1} + \Lambda \mathbf{A}^{-1} [(\nabla\Lambda) \mathbf{A} (\nabla \otimes \mathbf{A})] \mathbf{A}^{-1} \\ &\left. + [(\nabla\Lambda) \mathbf{A} (\nabla\Lambda)] \mathbf{A}^{-1} + \Lambda(\nabla\Lambda) \mathbf{A} \text{tr}_{2,3} [(\nabla \otimes \mathbf{A}) \mathbf{A}^{-1}] \mathbf{A}^{-1} \right\} \\ \mathbf{A} = \mathbf{1}: &\frac{1}{4\Lambda^2} \left\{ 3(\nabla\Lambda) \otimes (\nabla\Lambda) - 2\Lambda(\nabla \otimes \nabla\Lambda) + [(\nabla\Lambda)(\nabla\Lambda) - 2\Lambda(\Delta\Lambda)] \mathbf{1} \right\} \end{aligned}$$

$$\nabla T \implies \nabla \Lambda \implies \mathbf{D} \implies \sigma$$

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THANK YOU FOR YOUR ATTENTION!